# Simulation of Turbulent Premixed Hydrogen Combustion

J. Bell

Lawrence Berkeley National Laboratory

SIAM Conference on Computational Science and Engineering Miami, FL March 2-6, 2009

Collaborators: M. S. Day, V. E. Beckner, M. J. Lijewski

R. K. Cheng, S. Tachibana, M. Legrand



### Lean Premixed Turbulent Combustion



V-flame



4-jet Low-swirl burner (LSB)

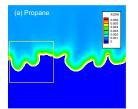


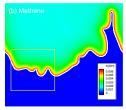
Slot burner

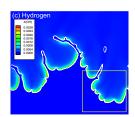
- Potential for efficient, low-emission power systems
- Design issues because of flame instabilities
- Focus on behavior of lean premixed hydrogen combustion



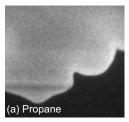
# Fuel dependence of flame structure

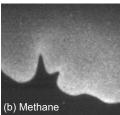


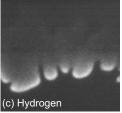




OH Mole fraction







**OH PLIF** 



### Compressible Navier Stokes

Gas phase combustion – mixture model for diffusion

$$\begin{array}{ll} \textbf{Mass} & \rho_t + \nabla \cdot \rho U = 0 \\ \textbf{Momentum} & (\rho U)_t + \nabla \cdot (\rho U U + \rho) = \rho \vec{g} + \nabla \cdot \tau \\ \textbf{Energy} & (\rho E)_t + \nabla \cdot (\rho U E + \rho U) = \nabla \cdot \kappa \nabla T + \nabla \cdot \tau U \\ & + \sum_m \nabla \cdot (\rho h_m D_m \nabla Y_m) \\ \textbf{Species} & (\rho Y_m)_t + \nabla \cdot (\rho U Y_m) = \nabla \cdot (\rho D_m \nabla Y_m) + \dot{\omega}_m \end{array}$$

#### Augmented with

- Thermodynamics
- Reaction kinetics
- Transport coefficients

Need to preserve chemical and transport fidelity



### Relevant Scales

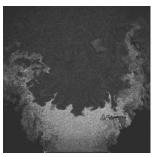
#### **Spatial Scales**

- Domain: ≈ 10 cm
- Flame thickness:  $\delta_T \approx$  1 mm
- Integral scale:  $\ell_t \approx 2-6 \text{ mm}$

#### Temporal Scales

- Flame speed  $O(10^2)$  cm/s
- Mean Flow: O(10<sup>3</sup>) cm/s
- Acoustic Speed: O(10<sup>5</sup>) cm/s

Fast chemical time scales but energy release coupling chemistry to fluid is on slower time scales



Mie Scattering Image



### Low swirl burner simulation

#### Simulation requirements

- No explicit model for turbulence, or turbulence/chemistry interactions
- Detailed chemistry based on fundamental reactions, detailed diffusion
- Simulate on time scales associated with the fluid velocity

Direct integration of compressible Navier Stokes too demanding

#### Exploit structure of the problem

- Mathematical model
- Approximation / discretization
- Solvers and software



### Mathematical formulation

Exploit natural separation of scales between fluid motion and acoustic wave propagation

Low Mach number model,  $M=U/c\ll 1$  (Rehm & Baum 1978, Majda & Sethian 1985)

Start with the compressible Navier-Stokes equations for multicomponent reacting flow, and expand in the Mach number, M = U/c.

Asymptotic analysis shows that:

$$p(\vec{x},t) = p_0(t) + \pi(\vec{x},t)$$
 where  $\pi/p_0 \sim \mathcal{O}(M^2)$ 

- $p_0$  does not affect local dynamics,  $\pi$  does not affect thermodynamics
- For open containers  $p_0$  is constant
- Pressure field is instanteously equilibrated removed acoustic wave propagation



### Low Mach number equations

$$\begin{split} & \text{Momentum} \quad \rho \frac{DU}{Dt} = -\nabla \pi + \nabla \cdot \left[ \mu \left( \frac{\partial U_i}{\partial x_j} + \frac{\partial U_j}{\partial x_i} - \frac{2}{3} \delta_{ij} \nabla \cdot U \right) \right] \\ & \text{Species} \quad \frac{\partial (\rho Y_m)}{\partial t} + \nabla \cdot (\rho U Y_m) = \nabla \cdot (\rho D_m \nabla Y_m) + \dot{\omega}_m \\ & \text{Mass} \quad \frac{\partial \rho}{\partial t} + \nabla \cdot (\rho U) = 0 \\ & \text{Energy} \quad \frac{\partial \rho h}{\partial t} + \nabla \cdot \left( \rho h \vec{U} \right) = \nabla \cdot (\lambda \nabla T) + \sum_m \nabla \cdot (\rho h_m D_m \nabla Y_m) \end{split}$$

Equation of state  $p_0 = \rho RT \sum_m \frac{Y_m}{W_m}$ 

System contains four evolution equations for  $U, Y_m, \rho, h$ , with a constraint given by the EOS.

Low Mach number system can be advanced at fluid time scale instead of acoustic time scale but . . .

We need effective integration techniques for this more complex formulation



### Constraint for reacting flows

Low Mach number system is a system of PDE's evolving subject to a constraint; differential algebraic equation (DAE) with index 3

Differentiate constraint to reduce index

$$\nabla \cdot U = \frac{1}{\rho c_p T} \left( \nabla \cdot (\lambda \nabla T) + \sum_m \rho D_m \nabla Y_m \cdot \nabla h_m \right) + \frac{1}{\rho} \sum_m \frac{W}{W_m} \nabla (D_m \rho \nabla Y_m) + \frac{1}{\rho} \sum_m \left( \frac{W}{W_m} - \frac{h_m(T)}{c_p T} \right) \dot{\omega}_m$$

Generalized projection method framework

- Finite amplitude density variation
- Inhomogeneous constraint
- Requires solution of variable coefficient, self-adjoint elliptic PDE



### Low Mach number numerics

### Fractional step scheme

- Advance velocity and thermodynamic variables
  - Advection
  - Diffusion
  - Stiff reactions
- Project solution back onto constraint variable coefficient elliptic PDE, multigrid
  Stiff kinetics relative to fluid dynamical time scales

$$\frac{\partial(\rho Y_m)}{\partial t} + \nabla \cdot (\rho U Y_m) = \nabla \cdot (\rho D_m \nabla Y_m) + \dot{\omega}_m$$

$$\frac{\partial(\rho h)}{\partial t} + \nabla \cdot (\rho U h) = \nabla \cdot (\lambda \nabla T) + \sum_{m} \nabla \cdot (\rho h_{m} D_{m} \nabla Y_{m})$$

#### Operator split approach

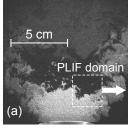
- Chemistry  $\Rightarrow \Delta t/2$
- Advection Diffusion  $\Rightarrow \Delta t$
- Chemistry  $\Rightarrow \Delta t/2$

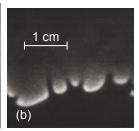
Coupled to block structured AMR



# Hydrogen combustion





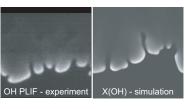


- OH PLIF shows gaps in the flame
- How do these flames burn?
- Are existing engineering models applicable?
- Can standard flame analysis techniques be used to analyze structure?

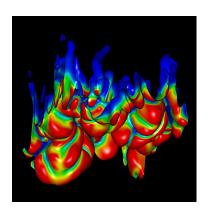


# Hydrogen flame in 3D

3D control simulation of detailed hydrogen flame at laboratory scales  $(3 \times 3 \times 9 \text{ cm domain}, \Delta x_f = 58 \mu\text{m})$ 



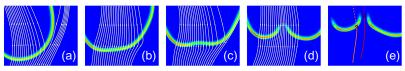
- Figure is "underside" (from fuel side of flame)
- Flame surface (isotherm) colored by local fuel consumption
- Cellular structures convex to fuel, robust extinction ridges



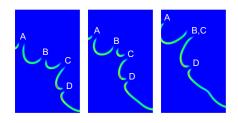


# Localized hydrogen flame "extinction"

#### Analysis from 2D study

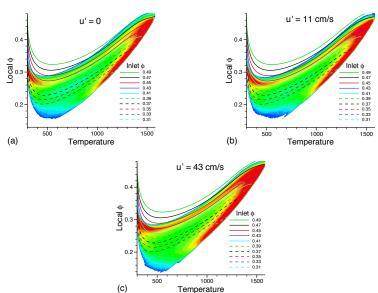


- Low-level localized strain event leads to onset of extinction.
- Lagrangian pathline analysis shows highly mobile fuel atoms diffuse "off-pathline", no fuel leakage.



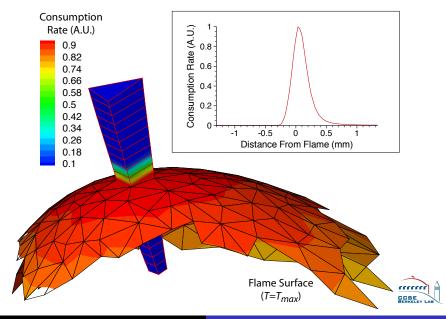


### Local flame enrichment

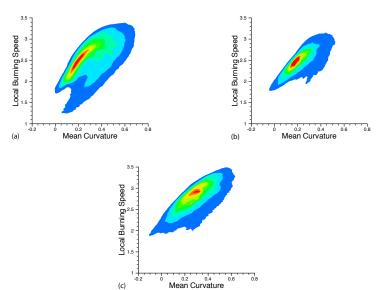




# Local consumption speed



# Flame speed versus curvature

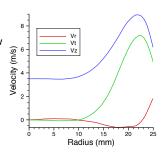




### Low swirl burner simulations

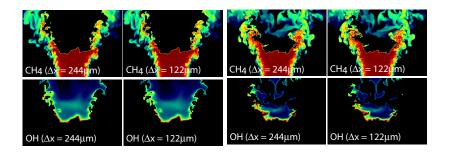
#### Strategy:

- Treat outflow from the nozzle as an inflow boundary condition
  - Mean flow and turbulent intensities from measured data
  - Impose synthetic turbulence as a perturbation to mean inflow
- Simulate flow in a rectilinear domain sitting above the outflow
- Four cases
  - Hydrogen ( $\phi = 0.37$ ) and methane ( $\phi = 0.7$ )
  - Laminar flame speed approximately 15 cm / sec
  - Two levels of mean flow and turbulence





### Methane swirl simulations

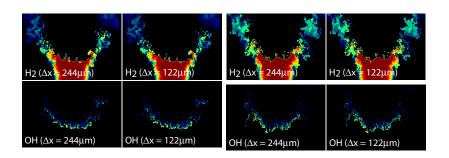


Weak Turbulence

Strong Turbulence



# Hydrogen swirl simulations

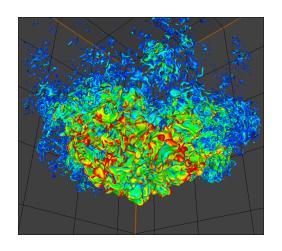


Weak Turbulence

Strong Turbulence

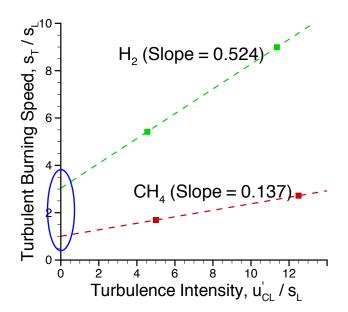


# Hydrogen flame surface





# Flame Speeds





# Summary

Developed new methodology to simulate realistic turbulent flames based on exploiting mathematical structure of combustion problems

- Range of scales relevant to laboratory experiments
- Detailed chemistry and transport
- No explicit models for turbulence or turbulence / chemistry interaction
- Methodology being applied to hydrogen flames in low-swirl burner

#### Future work

- Closed chamber simulations
- Include nitrogen chemistry for emissions
- High-pressure simulations

